

- $N=4$ susy Yang-Mills \rightarrow Hitchin eqns $\begin{cases} F_A - \phi \wedge \phi = 0 \\ d_A \phi = 0 \\ d_A * \phi = 0 \end{cases}$
 with target $C \times \mathbb{R}^2$
 \downarrow
 Riem. surface
 where A conn. on G -bundle
 $\phi \in \Omega^1(C, \text{ad } G)$

• Hitchin moduli space:

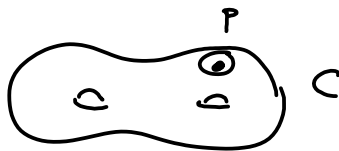
$$\begin{aligned} \mathcal{M}_H(G, C) &= \left\{ (A, \phi) \mid \begin{array}{l} F_A - \phi \wedge \phi = 0 \\ d_A \phi = 0 \\ d_A * \phi = 0 \end{array} \right\} / G \\ &= \left\{ \mathcal{A} = A + i\phi \mid \begin{array}{l} F_{\mathcal{A}} - \phi \wedge \phi = 0 \\ d_{\mathcal{A}} \phi = 0 \end{array} \right\} / G_{\mathbb{C}} \\ &= \mathcal{M}_{\text{flat}}^{G_{\mathbb{C}}}(C). \quad \text{Flat } G_{\mathbb{C}}\text{-connections} \end{aligned}$$

$\mathcal{M}_H(G, C)$ is hyperkähler. $\mathcal{I}, \mathcal{J}, \mathcal{K} = \mathcal{I}\mathcal{J}$
 $\omega_{\mathcal{I}}, \omega_{\mathcal{J}}, \omega_{\mathcal{K}}$

(\mathcal{J} is the "natural" \mathbb{C} str. in above description).

Then we have flexibility in choosing to consider either A or B model on $\mathcal{M}_H(G, C)$.

- Surface operators: $\begin{array}{l} \mathcal{M} = \mathbb{R}^2 \times C \\ \cup \\ \mathcal{D} = \mathbb{R}^2 \times p \end{array} \Rightarrow \text{Diagram}$



ie. consider a punctured Riem. surface (C, p) ,
 and allow (A, ϕ) to have a singularity at p .

Local model: on D' = punctured disc w/ coords $z = r e^{i\theta}$ ($p: z=0$).

$$\begin{cases} A = \alpha \frac{r d\theta}{r} + \dots \\ \phi = \beta \frac{dr}{r} - \gamma d\theta + \dots \end{cases} \quad \begin{array}{l} \alpha, \beta, \gamma \in \mathfrak{k} = \text{Lie}(\overline{T}) \\ \overline{T} \subset G \text{ max. torus} \end{array}$$

Gauge equivalence: $u \in \Lambda_{\text{Cochain}}(G)$, $f = e^{u\theta}$

acts by $\alpha \mapsto \alpha + u$

Hence: $(\alpha, \beta, \gamma) \in (\mathbb{T} \times \mathfrak{t} \times \mathfrak{t}) / \mathcal{W}$

- supersymmetry $N=4$ (parameters come in 4-tuples)
[e.g. in $N=2$ A-model, introduce B-field to complexify \mathcal{W}].

\Rightarrow need an extra "quantum" param. $\eta \in {}^L\mathbb{T}$ max. torus of Langlands dual of G .

[on $\mathcal{D} \cong \mathbb{R}^2$, add phase factor $\exp(i\eta \int_{\mathcal{D}} F_A)$]

$\Rightarrow (\alpha, \beta, \gamma, \eta) \in (\mathbb{T} \times \mathfrak{t} \times \mathfrak{t} \times {}^L\mathbb{T}) / \mathcal{W}$.

$\mathcal{M}_H(G, \mathbb{C}, p; \alpha, \beta, \gamma, \eta)$ $\dim H^2(\mathcal{M}_H(\dots)) = \text{rank } G$.

Model	Complex moduli	Kähler moduli
I	$\beta + i\gamma$	$\alpha + i\eta$
J	$\gamma + i\alpha$	$\beta + i\eta$
K	$\alpha + i\beta$	$\gamma + i\eta$

NB: B-model for J depends on $\gamma + i\alpha$ but not on \mathbb{C} str. of \mathbb{C}
(it appears on A-model for J)

However, e.g. B-model for I depends partially on \mathbb{C} str. of \mathbb{C} .

Ex: $G = SU(2)$, $\mathbb{C} = \mathbb{C} - \{0\}$ $\underbrace{\quad}_p$

$\Rightarrow \mathcal{M}_H \cong \mathbb{H}^2 // U(1)$ hK quotient at $\vec{\mu} = (\mu_I, \mu_J, \mu_K) = (\alpha, \beta, \gamma)$

for $\alpha = \beta = \gamma = 0$, get $\mathcal{M}_H \cong \mathbb{C}^2 / \mathbb{Z}_2$; otherwise $\mathcal{M}_H = T^*\mathbb{C}P^1$, with $\eta = \int_{\mathbb{C}P^1} \mathcal{B}$
(B-field for any of I, J, K)

can see e.g. by thinking of flat $G_{\mathbb{C}}$ connections

Ex: (Painlevé VI)

$$G = SU(2), \quad C = \mathbb{CP}^1 - \{p_1, p_2, p_3, p_4\}$$

$$\mathcal{M}_H(G, C, \dots)_T = \left\{ \rho: \pi_1(C) \rightarrow G_{\mathbb{C}} \mid \rho(\gamma_i) \in \mathcal{C}_i \right\} / \sim$$

conj. classes at the 4 punctures:
in terms of params. $V_i = \exp(2\pi i(\gamma_i + i\alpha_i))$

$$= \left\{ U_i \mid U_i \in \mathcal{C}_i, U_1 U_2 U_3 U_4 = \text{Id} \right\} / \sim.$$

Can write explicitly in terms of

$$\text{variables } \begin{cases} x_1 = \text{tr}(U_3 U_2) \in \mathbb{C} \\ x_2 = \text{tr}(U_1 U_3) \\ x_3 = \text{tr}(U_2 U_1) \end{cases} \quad \& \quad \text{params. } \begin{cases} \theta_1 = a_1 a_4 + a_2 a_3 \\ \theta_2 = \dots \\ \theta_3 = \dots \\ \theta_4 = a_1 a_2 a_3 a_4 + \sum_1^4 a_i^2 - 4 \end{cases}$$

where $a_i = \text{tr}(U_i) \quad i=1..4$
are determined by parameters α, γ .

$$\text{Then get } \mathcal{M}_H(\dots)_T = \left\{ (x_1, x_2, x_3) \in \mathbb{C}^3 \mid f(x_i, \theta_m) = 0 \right\}$$

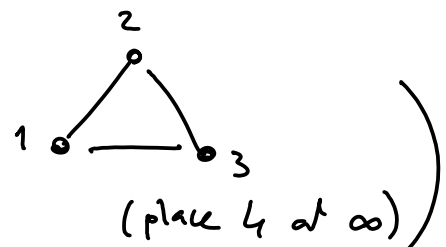
$$\text{where } f(x_i, \theta_m) = x_1 x_2 x_3 + \sum_1^3 (x_i^2 - \theta_i x_i) + \theta_4$$

\Rightarrow cubic surface in \mathbb{C}^3
w/ eqn depending on parameters.

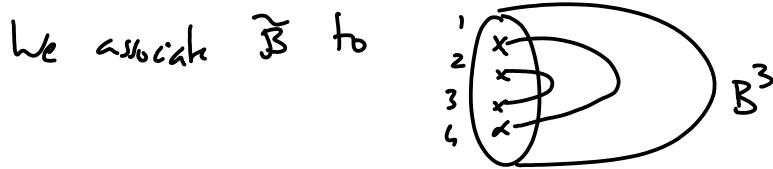
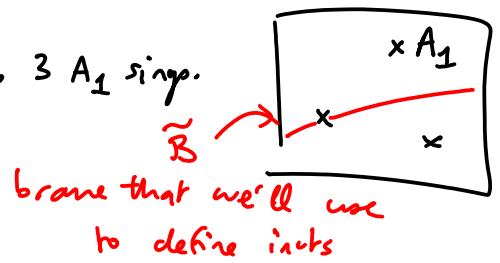
• $Br_3 \ni \sigma_i \quad i=1,2,3$ acts on $\mathcal{M}_H(\dots)_T$:

$$\begin{array}{ll} \sigma_i: & x_i \mapsto \theta_j - x_j - x_k x_i \quad \{i,j,k\} = \{1,2,3\} \\ & x_j \mapsto x_i \\ & x_k \mapsto x_k \end{array} \quad \begin{array}{l} \theta_i \mapsto \theta_j \\ \theta_j \mapsto \theta_i \\ \theta_k, \theta_4 \text{ stay} \end{array}$$

$$\left(\begin{array}{l} \sigma_i \sigma_j \sigma_i = \sigma_j \sigma_i \sigma_j \\ \sigma_k = \sigma_i \sigma_j \sigma_i^{-1} \end{array} \right)$$

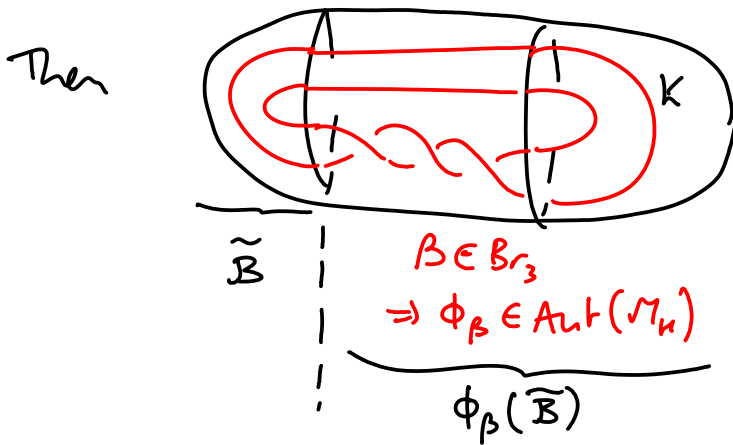


• Take $a_i = a$ all equal: then \mathcal{M}_n has 3 A_1 sing.



$$\begin{aligned} \text{ie. } \tilde{\mathcal{B}} &= \{ U_1 = U_4^{-1}, U_2 = U_3^{-1} \} \\ &= \{ x_1 = \text{tr}(U_3 U_2) = 2 \} \end{aligned}$$

$\tilde{\mathcal{B}}$ passes through one of the A_1 -singularities, but we define Hom 's ignoring the sing. pt.



$$\Rightarrow H_k := \text{Hom}(\tilde{\mathcal{B}}, \phi_\beta(\tilde{\mathcal{B}}))$$

Ex: $T_{2,k}$ torus knot: $\beta = \sigma_1^k$

$$\Rightarrow \dim H_{T_{2,k}} = 2\sigma(T_{2,k})$$